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GOALS, STRATEGIES AND FIRST DISCOVERIES OF AO327, THE ARECIBO ALL-SKY 327 MHz DRIFT PULSAR SURVEY

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ABSTRACT

We report initial results from AO327, a drift survey for pulsars with the Arecibo telescope at 327 MHz. The first phase of AO327 will cover the sky at declinations of -1° to 28° , excluding the region within 5° of the Galactic plane, where high scattering and dispersion make low-frequency surveys sub-optimal. We record data from a 57 MHz bandwidth with 1024 channels and $125 \mu\text{s}$ sampling time. The 60 s transit time through the AO327 beam means that the survey is sensitive to very tight relativistic binaries even with no acceleration searches. To date we have detected 44 known pulsars with periods ranging from 3 ms to 2.21 s and discovered 24 new pulsars. The new discoveries include 3 ms pulsars, three objects with periods of a few tens of milliseconds typical of young as well as mildly recycled pulsars, a nuller, and a rotating radio transient. Five of the new discoveries are in binary systems. The second phase of AO327 will cover the sky at declinations of 28° – 38° . We compare the sensitivity and search volume of AO327 to the Green Bank North Celestial Cap survey and the GBT350 drift survey, both of which operate at 350 MHz.

Key words: pulsars: general – stars: neutron

Online-only material: color figures

1. INTRODUCTION

Currently there are ~ 2000 known pulsars, the majority in our Galaxy and a few in the Magellanic Clouds. About 250 pulsars have rotation periods of less than 10 ms and are nature’s most accurate clocks. The unique properties of these millisecond pulsars (MSPs) can be exploited to study the exotic state of matter in the interiors of neutron stars, to detect gravitational waves, and to test the theory of general relativity in the strong gravity regime. For detecting gravitational waves using a Pulsar Timing Array (PTA; e.g., Jenet et al. 2005), one needs regular observations of MSPs that are extremely stable rotators and are, optimally, distributed all over the sky.

Most pulsars have been discovered in large radio surveys, but past survey coverage has favored the southern hemisphere and the region within a few degrees of the Galactic plane. The Arecibo drift-scan surveys carried out at 430 MHz prior to and during the Gregorian dome upgrade in the 1990s covered about 70% of the Arecibo sky (declinations between -0.5° and 38° ; Xilouris et al. 1998). Among the many pulsars found, some have been extremely important for science in general and radio astronomy in particular. These include: PSR B1534+12 (Stairs et al. 1998), the relativistic binary system which has allowed two different tests of general relativity; PSR B1257+12 (Wolszczan & Frail 1992), a MSP now known to be orbited by three planets and the first extra-solar planetary system discovered; and PSR J1713+0747 (Foster et al. 1993), a bright, stable MSP which dominates the sensitivity of the PTA (Demorest et al. 2013).

Despite the successes of the Arecibo searches made so far, they also had limitations. The sky was divided into five

regions, with each of these observed by a different group. Early observations were made with old equipment, which had coarser time and spectral resolution than is available from current instruments. Data collected during the upgrade were often of variable quality and the different groups were using different search algorithms and methods for interference excision. As a result, widely different success rates in finding new pulsars were reported. This complicates any efforts to draw inferences about the general pulsar population probed by these surveys. Given the patchy nature of the observations and data processing carried out so far, there is good reason to believe that $\sim 30\%$ of the Arecibo sky remains to be surveyed, and that many pulsars await discovery (Xilouris et al. 1998). This belief is supported by the discovery of 11 new pulsars in an analysis of 430 MHz drift data covering 490 deg^2 by Lorimer et al. (2004, 2005). Two of the new pulsars have periods < 60 ms.

For a few months in 2007, the Green Bank Telescope (GBT) underwent azimuth track repair, which made normal pointed observations impossible, but Lynch et al. (2013) and Boyles et al. (2013) used ~ 1500 hr to conduct a 350 MHz drift survey (GBT350) of sky in the declination range between -21° and 2° (which is mostly inaccessible to the Arecibo telescope) and to a lesser extent at declinations of 17° – 26° . GBT350 has discovered 35 pulsars, including 7 MSPs. One of the new MSPs is in a hierarchical triple system (S. M. Ransom et al. 2013, in preparation), one of only two such systems known. Another MSP, PSR J1023+0038, was the first object whose evolution from a low-mass X-ray binary to a radio-emitting MSP was observed almost in real time over the span of a decade (Archibald et al. 2009). A third discovery, the mildly recycled pulsar J0348+0432, is in a binary system with an orbital period of only 2.4 hr, and has the highest precisely measured neutron

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Figure 1. Current sky coverage of the AO327 survey. Gray strips indicate data taken. Yellow strips indicate data processed and searched. Black dots correspond to known pulsars from the ATNF catalog. Blue circles show known pulsars detected blindly by the AO327 search pipeline. Red circles show new pulsars discovered by AO327, and magenta circles show new pulsars discovered during the WAPP-based pilot 327 MHz drift survey from 2003.

(A color version of this figure is available in the online journal.)

star mass, $2.01 \pm 0.04 M_{\odot}$ (Antoniadis et al. 2013). This makes it a unique lab for testing alternative theories of gravity, because it probes the properties of gravity for the most extreme gravitational fields. Since 2009, the Green Bank North Celestial Cap survey (GBNCC; K. Stovall et al. 2013, in preparation) has searched the sky at declinations $>38^{\circ}$ (also inaccessible to the Arecibo telescope) at 350 MHz and has discovered 62 new pulsars, including 9 MSPs. The results of GBT350 and GBNCC demonstrate the excellent scientific potential of high-latitude pulsar surveys.

In Section 2 we discuss the current status of AO327 in terms of sky coverage, sensitivity, and search volume, and we make comparisons with the Green Bank low-frequency surveys. Section 3 focuses on the 24 new pulsars discovered by AO327 so far, and Section 4 presents our plans and projections for the next few years that will be necessary to survey the entire Arecibo sky at 327 MHz.

2. THE AO327 SURVEY

The one minute effective integration time of an Arecibo drift survey at 327 MHz means that it is sensitive to very tight relativistic binaries even with only moderate acceleration searches. The double pulsar J0737–3039 has an orbital period of 2.46 hr and was found in a Parkes survey with an integration time of 265 s without using an acceleration search during data processing (Burgay et al. 2006). This means that for similar flux densities, companion masses, and spin periods, AO327 is sensitive to comparable binaries with orbital periods of $2.46 \text{ hr} \times 60/265 = 0.57 \text{ hr}$ with a conventional periodicity search, and to even shorter-period comparable binaries if using acceleration search. Pulsars in tight relativistic binaries with orbital periods of a few hours or less can be timed very well at Arecibo, where the maximum source tracking time is $\sim 2.5\text{--}3 \text{ hr}$.

We conducted a pilot 327 MHz drift survey at Arecibo in 2003 for 100 hr. The frequency, lower than used in previous Arecibo drift surveys, was chosen because pulsars tend to have power-law spectra that decline steeply with increasing frequency

Table 1
Low-frequency Survey Parameters Used in Comparison

Parameter	AO327 Pilot	AO327	AO327	GBT350	GBNCC
Backend	WAPP	Mock	PUPPI	Spigot	GUPPI
f_{center} (MHz)	327	327	327	350	350
T_{rec} (K)	130	113	113	46	46
G (K Jy^{-1})	11	11	11	2	2
T_{obs} (s)	60	60	60	140	120
$\Delta\nu$ (MHz)	25	57	69	50	100
N_{ch}	512	1024	2816	2048	4096
Δt (μs)	256	125	82	82	82

($S \propto \nu^{\alpha}$, with average $\alpha = -1.7$). It also anticipated the installation of a new, cooled 327 MHz receiver, which was eventually commissioned in 2006. This pilot survey used the older Wideband Arecibo Pulsar Processor (WAPP; Dowd et al. 2000) backend and discovered three pulsars (Table 2, Figure 1).

With the cooled 327 MHz receiver and the Mock spectrometer⁷ available, we began taking data during a lengthy repair of the telescope platform in 2010, which restricted telescope movement for several months. The Mock spectrometer allows us to collect data from a 57 MHz bandwidth as opposed to 25 MHz for the WAPP-based pilot survey and has a better response to radio frequency interference (RFI) than the older backend.

In 2012, the Arecibo Observatory acquired a new backend, the Puerto Rico Ultimate Pulsar Processing Instrument (PUPPI), which is a clone of the similarly named GUPPI at Green Bank.⁸ PUPPI provides a “fast4k” search mode, allowing recording of 4096 channels with a minimum sampling time of 81.92 μs . After PUPPI control is integrated into the Arecibo telescope observing software, AO327 will begin using PUPPI for regular survey observations. In the meantime, we have successfully used PUPPI in this mode to confirm some of the new pulsars in Table 2, and in coherent dedispersion mode we will use it for timing AO327 discoveries. Because of hardware constraints, the smallest bandwidth available with PUPPI is 100 MHz, while the 327 MHz receiver bandwidth is 67 MHz. A modification in the PUPPI data-taking software will allow us to record data only from the 68 MHz receiver bandwidth although the 4096 channels of the “fast4k” mode cover 100 MHz. This accounts for the 2816 channels listed in Table 1 for the future AO327 PUPPI setup.

2.1. Coverage

The AO327 survey aims to search the entire Arecibo sky (declinations between -1° and 38°) for pulsars and radio transients at 327 MHz. We plan to do this in two phases. During the 2010 platform repair, telescope movement was restricted successively to several different azimuth and zenith angle ranges, giving access to declinations (decl.) of $0^{\circ}\text{--}28^{\circ}$. The necessity to coordinate observations with repair work and regularly scheduled projects that were able to use the telescope in this restricted mode resulted in patchy sky coverage. After the movement restrictions were lifted in early 2011, AO327 continued to observe during time slots that were unassigned or became available at short notice (e.g., due to radar transmitter malfunction) until a proposal for dedicated telescope time was submitted and accepted.

Our Phase 1 goal is to fill in coverage gaps in the decl. = $0^{\circ}\text{--}28^{\circ}$ range (Figure 1), excluding the region within $\pm 5^{\circ}$ of the

⁷ <http://www.naic.edu/~phil/hardware/pdev/pdev.html>

⁸ <http://safe.nrao.edu/wiki/bin/view/CICADA/GUPPISupportGuide>

inner Galactic plane. The PALFA survey (Cordes et al. 2006) already covers the latter region at 1.4 GHz, a frequency better suited for pulsar searching in the high dispersion and scattering environment of the Galactic plane. We also include in Phase 1 the remaining 1° strip at the south edge of sky accessible to Arecibo (decl. = -1° – 0°) that was not covered during the period of platform repair. Phase 2 will extend the survey to the remaining portion of the Arecibo sky, decl. = 28° – 38° . So far the AO327 has taken survey data with the Mock spectrometer for a total of 1687 hr. In order to completely cover the field shown in Figure 1, a total of 2707 hr are necessary. So far we have 1413 hr of non-redundant sky coverage and therefore at least 1294 hr are necessary to fill in all gaps in the Phase 1 survey region. AO327 gets ~ 400 hr of telescope time per year. However, this includes setup and slewing overhead in addition to actual data-taking time. In addition, some observing time is used for candidate confirmations. We estimate that Phase 1 will be completed by 2017. Phase 2 of the survey will take two additional years.

The 274 hr of data corresponding to redundant sky coverage are due to limited telescope mobility in 2010 and the combined partial overlap in right ascension (R.A.) of adjacent drift strips. In addition, the $15'$ strip of sky centered at decl. = $19^\circ 52'$ has been observed for 2.5×24 hr. This decl. corresponds to the telescope stow position. During observatory shut-down due to tropical storm or hurricane warnings as well as in actual severe tropical weather conditions, the telescope is locked in the stow position, but it is still possible to take data in drift mode. AO327 observes under such conditions and we have detected several known pulsars in data from this “hurricane strip.” Redundant sky coverage improves the chance of detecting transient, repeating sources, scintillating, eclipsing, and intermittent pulsars, and weak or long-period pulsars whose detectability in a blind search may be affected by variable RFI. Case in point is PSR J0241+16 (Table 2), which was discovered in the pilot WAPP-based 327 MHz drift search from 2003 and has been confirmed in a longer observation, but was not detected via a blind search in Mock AO327 data covering the same strip of sky. Folding the corresponding data span with the pulsar’s period (P) and dispersion measure (DM) produced a marginal detection contaminated by RFI. This pulsar also has the second lowest DM among our discoveries, which makes it especially susceptible to scintillation.

2.2. Sensitivity

We calculate the minimum detectable pulsar flux density for a survey using the equation

$$S_{\min} = \frac{\beta (S/N)_{\min} T_{\text{sys}}}{G \sqrt{N_p T_{\text{obs}} \Delta \nu}} \sqrt{\frac{W}{P - W}}, \quad (1)$$

where $(S/N)_{\min} = 6$ is the minimum signal-to-noise ratio corresponding to a detection, $T_{\text{sys}} = T_{\text{rec}} + T_{\text{sky}}$ is a system temperature consisting of receiver and sky terms, G is the gain, $N_p = 2$ is the number of polarization channels summed, T_{obs} is the maximum transit time through the beam, $\Delta \nu$ is the bandwidth, W is the observed pulse width, and P is the pulsar period. The observed pulse width $W = \sqrt{W_{\text{int}}^2 + \Delta t_{\text{DM}, 1\text{ch}}^2 + \Delta t_{\delta\text{DM}}^2 + \Delta t_{\text{samp}}^2 + \tau_s^2}$. For our calculations, we use $T_{\text{sky}} = 50$ K (Haslam et al. 1982) and assume an intrinsic pulse width $W_{\text{int}} = 0.1P$. The factor β expresses sensitivity degradation due to digitization effects. It is unity for the Mock spectrometer, PUPPI, and GUPPI, and ≈ 1.16 for three-level correlators like the Spigot and WAPP.

Table 2
New Pulsars Discovered by the AO327 Drift Survey

Name	R.A. (hh:mm:ss) ^a	Decl. (dd:mm)	P (ms)	DM (pc cm ⁻³)	SP	Notes
J0158+21	01:58:29	21:08	505.3	20		
J0229+20	02:29:03	20:58	806.9	27	Y	Nuller
J0241+16 ^b	02:41:46	16:04	1545.4	16	Y	
J0244+14	02:44:51	14:27	2128.1	31	Y	
J0453+16	04:53:38(7)	16:01(1.7)	45.8	30		Binary
J0457+23	04:57:06(7)	23:34(1.7)	504.9	59	Y	
J0509+08	05:09:24(7)	08:57(1.7)	4.1	38		Binary
J0608+00	06:08:49	00:39	1076.2	48	Y	
J0628+06	06:28:04	06:38	346.5	87		
J0806+08	08:06:05	08:17	2063.1	46	Y	
J0824+00	08:24:16	00:27	9.9	35		Binary
J0848+16 ^b	08:48:53	16:43	452.4	38	Y	
J0928+06	09:28:44	06:14	2060.4	50	Y	
J1010+15 ^b	10:10:00	15:51	...	42	Y	RRAT
J1726-00	17:26:23	-00:15	1308.6	57		
J1802+03	18:02:44	03:38	664.3	77		
J1807+04	18:07:25	04:05	798.9	53	Y	
J1821+01	18:21:48(7)	01:52(1.7)	33.8	52		
J1937-00	19:37:09	-00:17	240.1	68.6		
J1945+07	19:45:55	07:17	1073.9	62	Y	
J2204+27	22:04:40(7)	27:02(1.7)	84.7	35		Binary
J2234+06	22:34:22(7)	06:11(1.7)	3.6	11		Binary
J2329+16	23:29:50	16:57	632.1	31	Y	
J2340+08	23:40:45(7)	08:33(1.7)	303.3	24	Y	

Notes. All pulsars except PSR J1010+15 were discovered via periodicity search. The column labeled “SP” shows whether an object was also detected in a single-pulse search.

^a R.A. and decl. are given in the J2000 coordinate system. The uncertainties in both coordinates are $7/5$, the 327 MHz beam radius, unless otherwise indicated. The positions with smaller uncertainties were obtained from follow-up observations at 1.4 GHz.

^b Found in AO327 Pilot data from 2003.

To calculate the dispersion smearing over one channel ($\Delta t_{\text{DM}, 1\text{ch}}$) and the dispersion smearing due to DM error in the worst case, when the pulsar DM is halfway between two trial DMs ($\Delta t_{\delta\text{DM}}$), we use the cold plasma dispersion relation (e.g., Lorimer & Kramer 2005). The term Δt_{samp} is equal to the actual sampling time for low DMs, and reflects downsampling for searching higher DMs, where the scattering broadening τ_s dominates. We calculate τ_s using the empirical fit of observed scattering versus DM from Bhat et al. (2004).

In order to compare the boresight sensitivity of several low-frequency, high-latitude surveys, we use the parameters in Table 1. We scale the GBT350 and GBNCC sensitivity to 327 MHz using a spectral index $\alpha = -1.7$. The current Mock spectrometer-based version of AO327 is significantly more sensitive than GBT350 or GBNCC to MSPs at low DMs and matches GBT350 sensitivity at DM ~ 50 pc cm⁻³ (Figure 2). With the new PUPPI backend, AO327 will be more sensitive to MSPs at any DM than either GBT350 or GBNCC. We expect new MSP discoveries to have relatively low DMs, as according to Bhat et al. (2004), for a DM of 100 pc cm⁻³ and $f_{\text{center}} = 327$ MHz, $\tau_s = 1$ ms. The detectability of high-DM MSPs in low-frequency surveys is limited by scattering.

Another important quantity to consider is the cross-over point between a dispersion-limited and a scattering-limited detection regime as a function of trial DM. In the dispersion-limited regime, we can gain sensitivity to fast pulsars by minimizing the correctable broadening due to the sampling time,

Figure 2. Minimum detectable flux density at beam center vs. pulse period for the past (WAPP), present (Mock), and future (PUPPI) incarnations of AO327, as well as the GBT350, and GBNCC surveys for three DMs typical of pulsars at high Galactic latitudes. The sensitivity of GBT350 and GBNCC was scaled to 327 MHz using a spectral index $\alpha = -1.7$.

(A color version of this figure is available in the online journal.)

downsampling of the data during processing, and dispersion broadening caused by the finite width of a frequency channel and the finite spacing between successive DMs in the trial DM list, $\Delta t = \sqrt{\Delta t_{\text{DM,1ch}}^2 + \Delta t_{\text{DM}}^2 + \Delta t_{\text{samp}}^2}$. In the scattering-limited regime, uncorrectable scattering broadening limits the detectability of fast pulsars and improvements in the data-taking setup or hardware give diminishing returns. Figure 3 shows this cross-over point for the surveys in Table 1. Ideally, the smallest practical DM spacing and no downsampling should be used for search DMs smaller than the cross-over DM. In practice, some sensitivity is sacrificed in favor of shorter data processing times. For search DMs larger than the cross-over DM, downsampling factors and DM spacings are chosen such that $\Delta t \ll \tau_s$.

2.3. Search Volume

A key question to consider is how much volume is searched in a survey. The search volume is a strong function of frequency, direction and distance and is affected by propagation effects (dispersion and scattering), search sensitivity $S_{\min}$ and solid angle Ω . It is this volume that largely determines the total number of objects detectable, N . Assuming that the survey depth does not reach the edge of the pulsar population, but a smaller distance $d_{\max}$, then $N \propto V_{\max} = (1/3)\Omega d_{\max}^3$. Here, $V_{\max}$ is the maximum spatial volume probed. For a given pulsar at distance D with flux density S , $V_{\max} \propto D^3(S/S_{\min})^{3/2}$. Comparing AO327 and GBT350, the ratio of survey volumes is

$$\frac{V_{\max, \text{AO327}}}{V_{\max, \text{GBT350}}} = \left[\frac{S_{327\text{MHz}}/S_{350\text{MHz}}}{S_{\min, \text{AO327}}/S_{\min, \text{GBT350}}} \right]^{3/2}. \quad (2)$$

For pulsars with median spectral index $\alpha = -1.7$, the numerator $S_{327\text{MHz}}/S_{350\text{MHz}} = 1.12$. Figure 4 shows $V_{\max, \text{AO327, Mock}}/V_{\max, \text{GBT350}}$ and $V_{\max, \text{AO327, Mock}}/V_{\max, \text{GBNCC}}$ versus pulse period. For long-period pulsars, whose detectability is not affected by scattering and dispersion, AO327_{Mock} is searching a four times larger volume per unit solid angle than

Figure 3. Illustration of the cross-over point between a dispersion-limited and scattering-limited detection regime for the surveys we compare in this paper. In both panels, a thin solid black line shows the uncorrectable scattering broadening τ_s as a function of DM according to Bhat et al. (2004). Thin dashed black lines show $\pm$ one order of magnitude from τ_s , which is approximately the scatter around this empirical relation. In the top panel, a blue line shows the correctable broadening $\Delta t = \sqrt{\Delta t_{\text{DM,1ch}}^2 + \Delta t_{\text{DM}}^2 + \Delta t_{\text{samp}}^2}$ for GBT350, and a green line shows the same quantity for GBNCC. In the bottom panel, Δt is denoted by a thick black line for the pilot AO327 survey, and by red and cyan lines for AO327_{Mock} and AO327_{PUPPI}, respectively. For trial DMs lower than the DM where $\Delta t(\text{DM})$ and $\tau_s(\text{DM})$ cross for a survey, the survey is dispersion-limited. For trial DMs above that point, it is scattering-limited. The plot takes into account downsampling factors and DM spacings, which increase discretely as trial DMs increase. This accounts for the steps in Δt . AO327_{PUPPI} matches closely the cross-over point and the $\Delta t(\text{DM})$ dependence of GBNCC, the more sensitive of the GBT surveys, while having the advantage of a lower frequency, where pulsars are brighter.

(A color version of this figure is available in the online journal.)

Figure 4. Ratios of maximum volume per unit solid angle searched (Equation (2)) for three DMs typical of pulsars at high Galactic latitudes. Solid lines denote $V_{\max, \text{AO327-Mock}}/V_{\max, \text{GBNCC}}$ and dashed lines denote $V_{\max, \text{AO327-Mock}}/V_{\max, \text{GBT350}}$ for the Mock spectrometer setup currently used for AO327 observations.

(A color version of this figure is available in the online journal.)

GBT350 and a 2.7 times larger volume per unit solid angle than GBNCC. For DMs typical of expected high-latitude MSP discoveries, AO327_{Mock} is searching ~ 0.7 –4 times the volume of GBT350 and ~ 0.4 –1.8 times the volume of GBNCC. The advantages of the GBT surveys in the case of detecting very fast, higher DM pulsars come from their smaller sampling times and

Figure 5. Ratios of maximum volume per unit solid angle searched (Equation (2)) for three DMs typical of pulsars at high Galactic latitudes. Solid lines denote $V_{\max, \text{AO327-PUPPI}} / V_{\max, \text{GBNCC}}$ and dashed lines denote $V_{\max, \text{AO327-PUPPI}} / V_{\max, \text{GBT350}}$ for the new PUPPI backend that AO327 will transition to in the second half of 2013.

(A color version of this figure is available in the online journal.)

narrower channels compared to AO327_{Mock} (Table 1), as well as the slightly less pronounced scattering and dispersion broadening at 350 MHz compared to 327 MHz. However, that advantage will be drastically reversed in favor of AO327 once we start using the new PUPPI backend (Figure 5). AO327_{PUPPI} will search a 3.1 times larger volume per unit solid angle than GBNCC and a 4.7 times larger volume than GBT350 for slow pulsars. For MSPs, AO327_{PUPPI} will search ~ 2.6 – 3.3 times the volume of GBNCC and ~ 4.8 – 7.2 times the volume of GBT350.

2.4. Data Processing

Data are recorded by the Mock spectrometer as 16 bit short integers in the PSRFITS data format. We have determined that 4 bits per data point provide enough dynamic range so that sensitivity is unaffected, and we routinely rescale and rewrite data files in a 4 bit PSRFITS format for shipping to processing sites and for long-term storage in order to minimize disk needs. Data are processed at the University of Texas at Brownsville (UTB), West Virginia University (WVU), and the Arecibo Observatory (AO) using a pipeline based on the PRESTO pulsar search code.⁹ The AO beam width at 327 MHz is $15'$, giving a maximum transit time of 1 minute. Data are processed in 60 s spans, overlapping successive spans by 30 s. Data are dedispersed with 6358 trial DMs in the range 0–1000 pc cm^{-3} and searched for periodic and transient signals. We defer a more detailed and technical discussion of the search pipeline to a future paper.

Periodic and single-pulse candidates are stored in databases at UTB and WVU. A candidate ranking method based on the PEACE algorithm (Lee et al. 2013) assigns a grade to each periodic candidate corresponding to the candidate’s likelihood of being a real pulsar. The grade determines in what order candidates are examined by participants from AO, UTB, WVU, and the University of Wisconsin at Milwaukee via web portals hosted at UTB and WVU. There are currently more than 50

Figure 6. Barycentric period measurements of PSR J2234+06 vs. epoch. Black dots denote observations and the red curve shows the expected change in observed period due to eccentric orbital motion for the best-fit orbit. In this case, the orbital period is 32 days, and $e = 0.13$.

(A color version of this figure is available in the online journal.)

AO327 participants with various levels of involvement in the survey and expertise, ranging from high school students to university professors.

3. FIRST DISCOVERIES

So far the AO327 search pipeline has detected 44 known pulsars with periods in the range 0.003–2.21 s and discovered 24 pulsars (Table 2), including 3 MSPs, 3 objects with periods typical of young as well as mildly recycled pulsars, a nuller, and a rotating radio transient (RRAT). Five of the pulsars in Table 2 are included in an Arecibo timing campaign that began in 2013 February. We expect to begin timing the rest of the new discoveries in mid- to late 2013. Four of the new pulsars are within 10° of the Galactic plane. Two of these pulsars are close to the portion of the inner Galaxy visible from Arecibo, and one, PSR J1821+01, has $P \sim 34$ ms. PALFA searches the region within 5° of the Galactic plane in the portion of the Galaxy accessible to the Arecibo telescope: $32^\circ < l < 77^\circ$ and to a lesser extent $168^\circ < l < 214^\circ$. The fact that a low-frequency survey like AO327 is finding fast pulsars close to the inner Galactic plane underscores the necessity for mid-latitude targeted pulsar searches in this part of the sky. Ideally future surveys of the $5^\circ \lesssim |b| \lesssim 15^\circ$ region will use frequencies higher than 327 MHz in order to minimize scattering and make it possible to probe deeper parallel to the Galactic plane, but lower than the PALFA frequency of 1.4 GHz, where pulsars are significantly weaker and the number of pointings required would be prohibitively large.

3.1. PSR J2234+06

PSR J2234+06 is an MSP with $P = 3.58$ ms and a 32 day orbit (Figure 6). Assuming a pulsar mass of $1.4 M_\odot$, the companion star has a minimum and median mass of about 0.19 and 0.23 $M_\odot$, respectively. The orbital eccentricity of this system, $e = 0.13$, is much higher than observed for most other Galactic MSPs, which have orbits with eccentricities on the order of 10^{-7} to 10^{-3} . This trend is generally understood to be an inevitable consequence of orbital circularization during the accretion episode that leads to the recycling of the pulsar (e.g.,

⁹ <http://www.cv.nrao.edu/~sransom/presto>

Phinney & Kulkarni 1994). The spin and orbital parameters of PSR J2234+0611 make it a more compact, less massive cousin to PSR J1903+0327, an MSP with $P = 2.15$ ms, $P_b = 95$ days, $e = 0.44$, and a main sequence companion with a mass of $\sim 1 M_\odot$ (Champion et al. 2008). Such systems cannot be formed within the current understanding of stellar evolution and are believed to originate either from an exchange interaction in a region with a high stellar density or from a hierarchical triple that becomes dynamically unstable due to mass transfer to the pulsar from a third system component that is eventually ejected. Freire et al. (2011) reach the latter conclusion about PSR J1903+0327 based on the motion of the system in the Galaxy and Portegies-Zwart et al. (2011) independently reach the same conclusion based on numerical simulations of the evolution of the system.

Follow-up observations of PSR J2234+06 at 1.4 GHz revealed that the pulsar is very bright and can be easily timed at this frequency. Based on the folded pulse profiles at 1.4 GHz and 327 MHz, we estimate a peak flux density of 15 mJy at 1.4 GHz and 4.4 mJy at 327 MHz. The corresponding period-averaged flux densities are 0.7 and 0.3 mJy, respectively. This suggests that the pulsar either scintillates strongly on a timescale of minutes or has an anomalous spectrum rising with frequency. Most pulsars exhibit power-law spectra, $S \propto \nu^\alpha$, with an average spectral index $\alpha = -1.7$, and MSPs spectra on average decline even more steeply with increasing frequency. PSR J2234+06 has the lowest DM of all AO327 discoveries so far, which makes it especially prone to scintillation. One 1.4 GHz timing observation (2013 March 25) did not produce a detection. We rule out instrumental problems as other pulsars observed with the same setup before and after PSR J2234+06 were successfully detected in the same session. The pulsar was again detected easily at 1.4 GHz in subsequent timing observations. Apart from scintillation, another possible explanation of the one non-detection is that PSR J2234+06 is in an eclipsing binary system.

Time-of-arrival (TOA) measurements extracted from our 1.4 GHz observations have uncertainties on the order of only $0.4 \mu\text{s}$ within 2.5 minutes. Since $\sigma_{\text{TOA}} \propto 1/\sqrt{T_{\text{obs}}}$, we expect TOAs of this pulsar to have uncertainties on the order of $0.11 \mu\text{s}$ within half an hour, comparable to the typical precisions achieved within a half-hour observation for pulsars included in the PTAs now being used to search for gravitational waves (Demorest et al. 2013; Hobbs et al. 2010).

The brightness of PSR J2234+06 at 1.4 GHz and the orbital eccentricity imply that, as in the case of PSR J1903+0327, we will certainly be able to measure the rate of advance of periastron ($\dot{\omega}$). If the companion is compact, $\dot{\omega}$ is solely due to the effects of general relativity and can be used to estimate the total mass of the system (Weisberg & Taylor 1981). If the orbit is highly inclined with respect to the plane of the sky, the excellent timing precision of PSR J2234+06 means we will likely be able to measure the Shapiro delay, and therefore obtain a mass ratio and determine the masses of both components.

3.2. PSR J0453+16

PSR J0453+16 has $P = 45.8$ ms and $\text{DM} = 30 \text{ pc cm}^{-3}$. The discovery observation was made on 2010 May 30 and the pulsar was confirmed on 2012 August 21. The effective barycentric period derivative between these two observations is $\dot{P}_{\text{eff}} = -1.47(7) \times 10^{-13} \text{ s s}^{-1}$. Several subsequent observations show sign changes in $\dot{P}_{\text{eff}}$. Since rotation-powered pulsars typically spin down as energy is transferred to the emission mechanism, we conclude that the change in period observed for this pulsar is affected by orbital motion in a binary system.

Figure 7. Barycentric period measurements of PSR J0453+16 vs. epoch. Black dots denote observations and the red curve shows the expected change in observed period due to eccentric orbital motion for the best-fit orbit. In this case, the orbital period is 4 days, and $e = 0.11$.

(A color version of this figure is available in the online journal.)

A fit for the orbit (Figure 7) indicates an orbital period of four days and an eccentricity of 0.11. The projected semi-major axis is 14.5 lt-s . Using these parameters and assuming a pulsar mass of $1.4 M_\odot$, we get a minimum and mean companion masses of $1.0 M_\odot$ and $1.3 M_\odot$, respectively. Given its high eccentricity, mean companion mass of about $1.3 M_\odot$, and distance from the Galactic plane ($\sim 0.32 \text{ kpc}$), PSR J0453+16 may be in a double neutron star system. This pulsar is also easily detectable at 1.4 GHz and along with the orbital eccentricity this implies that, as the case of PSR J2234+06, we will be able to measure $\dot{\omega}$ and estimate the total mass of the system.

3.3. PSR J0509+08

PSR J0509+08 has $P = 4.06$ ms and $\text{DM} = 38 \text{ pc cm}^{-3}$. The pulsar was discovered in an observation on 2010 May 2 and confirmed on 2013 April 19. We use the barycentric periods and respective uncertainties reported by Presto for these data sets to calculate the effective barycentric period derivative between the two observations, $\dot{P}_{\text{eff}} = 2.83(3) \times 10^{-15} \text{ s s}^{-1}$. MSPs typically have period derivatives on the order of 10^{-20} – $10^{-19} \text{ s s}^{-1}$. Of 151 pulsars with $P < 0.015 \text{ s}$ and measured $\dot{P}$ in the ATNF catalog, only 2 have $\dot{P} \sim 10^{-18} \text{ s s}^{-1}$ and none have a larger $\dot{P}$. We therefore conclude that the observed period change for PSR J0509+08 is due to its orbital motion in a binary system. We estimate the flux density of this pulsar to be 5 mJy at 327 MHz and 0.8 mJy at 1.4 GHz, corresponding to a spectral index of -1.3 . A TOA measurement extracted from a 20 minute follow-up observation at 327 MHz has uncertainty of only $0.2 \mu\text{s}$, indicating that PSR J0509+08, like PSR J2234+06, is an excellent candidate for inclusion in PTAs. A TOA extracted from a four minute observation at 1.4 GHz has uncertainty of $3.7 \mu\text{s}$, corresponding to a projected TOA uncertainty of $1.3 \mu\text{s}$ within half an hour. We expect the timing precision at 1.4 GHz, a frequency typically used for PTA timing, to improve after we obtain a more precise position from regular timing observations and can observe the pulsar at boresight.

3.4. PSR J0824+00

PSR J0824+00 has $P = 9.86$ ms and $DM = 35$ pc cm $^{-3}$. For this pulsar, $\dot{P}_{\text{eff}} = 2.4(1) \times 10^{-14}$ s s $^{-1}$ between the discovery observation from 2012 September 2 and the confirmation from 2013 April 16. This suggests orbital motion for the same reasons as for PSR J0509+08.

3.5. PSR J2204+27

PSR J2204+27 has $P = 84.7$ ms and $DM = 35$ pc cm $^{-3}$. $\dot{P}_{\text{eff}}$ between the discovery observation on 2010 May 22 and the confirmation on 2011 August 31 is $1.8(3) \times 10^{-13}$ s s $^{-1}$. Five observations at 327 MHz within 24 days in 2011 September do not show a change in the observed period to within the precision reported by Presto. However, the effective period derivative between 2011 September and the start of our timing campaign in 2013 February is $-4.7(9) \times 10^{-14}$ s s $^{-1}$, again pointing to binary motion as the explanation. Since no $\dot{P}_{\text{eff}}$ was observed on a timescale of days, this suggests that the orbital period is on the order of months or longer.

3.6. PSR J1821+01

PSR J1821+01 has $P = 33.8$ ms and $DM = 52$ pc cm $^{-3}$. The calculated uncertainty in $\dot{P}_{\text{eff}}$ is 1.1×10^{-13} s s $^{-1}$ between the 2012 September 24 discovery and the 2012 November 3 confirmation, and 1.1×10^{-15} s s $^{-1}$ between the confirmation and our first timing observation from 2013 April 18. We did not detect any $\dot{P}_{\text{eff}}$ significantly larger in magnitude than these uncertainties. The period of PSR J1821+01 is consistent with either a young or a mildly recycled pulsar. Young pulsars have $\dot{P} \sim 10^{-14}$ – 10^{-12} s s $^{-1}$, and mildly recycled pulsars have $\dot{P} \sim 10^{-17}$ – 10^{-19} s s $^{-1}$. Our lack of detection of a $\dot{P}_{\text{eff}}$ even between observations separated by 166 days suggests that PSR J1821+01 is an isolated mildly recycled pulsar. PSR J1821+01 has $b \approx 7^\circ$, and the catalog of Galactic supernova remnants (Green 2009) has no nearby entries. The X-ray source 1RXS J182217.2+014227 is at the edge of the discovery beam, ~ 7.5 from beam center. We ruled out association with the pulsar by pointing at the X-ray source immediately after the successful confirmation observation and observing it with the same 327 MHz telescope setup and 10 minute integration time. This pointing produced no detection.

3.7. PSR J0229+20

PSR J0229+20 has $P = 806.9$ ms and $DM = 27$ pc cm $^{-3}$, and was discovered via a periodicity search. Figure 8 shows the pulse profile averaged over the ~ 60 s of the discovery observation, as well as a grayscale panel of subintegration versus pulse phase. The length of a subintegration for the purpose of plotting is partially determined by the period of a pulsar candidate, and in this case each row in the grayscale panel corresponds to an average of ~ 4 s of data. Darker regions indicate stronger emission. The observed intensity variations between subintegrations may be affected by intrinsic variability and/or scintillation. The emission in the first subintegration ($t = 0$ – 4 s) of the data span shown is weak as the pulsar enters the half-power contour of the telescope beam. The subintegration at $t \approx 24$ – 28 s shows little to no emission compared to other subintegrations. During the last subintegration ($t \gtrsim 60$ s) the pulsar begins to exit the telescope beam and emission intensity decreases. The beam of the 327 MHz Arecibo receiver is elliptical ($\approx 15' \times 14'$). In addition, the parallactic angle varies

Figure 8. Folded pulse profile (top) and pulse phase vs. time (bottom) for the discovery observation of PSR J0229+20. A gap in emission is evident at $t \sim 25$ s and may be due to nulling or scintillation.

between drift tracks done at different azimuth and zenith angles. These two properties account for the fact that even for steady emitters, the observed intensity versus time generally does not trace out an ideal Gaussian beam cross-section.

Figure 9 shows a plot of the same data span produced by the single-pulse branch of our search pipeline. The top three panels show statistics for bursts with $S/N > 5$. The prominent peaks in the number of pulses and S/N versus trial DM plots are indicative of a confident detection of many bright dispersed pulses. The observed single-pulse S/N and total pulse count are highest at the trial DM closest to the actual pulsar DM. In the bottom panel, pulse intensity gradually increases at the start of the observation as the pulsar enters the beam. However, there are no obvious non-background events at $t \approx 22$ – 28 s, corresponding to the gap in periodic emission from Figure 8. Nulling pulsars typically shut off abruptly, and just as abruptly start emitting again. This is consistent with bright single pulses bracketing the absence of detections at $t \approx 22$ – 28 s.

We calculate the diffractive scintillation timescale Δt_{DISS} for PSR J0229+20 based on equations in Section 4 of Lorimer &

Figure 9. Single-pulse search output for the discovery observation of PSR J0229+20. The top panels show the number of pulses vs. S/N and DM, and S/N vs. DM for events with S/N > 5. At bottom the same events are plotted vs. DM and time. Brighter events are represented by larger circles. Events are weaker at the start of the data span shown as the pulsar enters the telescope beam. In contrast, the gap at $t \sim 22\text{--}28$ s has no single-pulse events and is flanked by bright pulses before and after. This may be due to nulling or scintillation.

Kramer (2005), the average two-dimensional speed for non-recycled pulsars of $\sim 246 \text{ km s}^{-1}$ (Hobbs et al. 2005), and a distance estimate of 1.1 kpc based on the pulsar DM and position, and the NE2001 model of ionized gas distribution in the Galaxy (Cordes & Lazio 2002). We obtain $\Delta t_{\text{DISS}} = 54$ s. However, the empirical relation from Bhat et al. (2004) used to calculate the scattering broadening time has a scatter of an order of magnitude, which makes Δt_{DISS} values of 17–173 s plausible, and we cannot completely rule out scintillation as the cause of the observed brightness variability of PSR J0229+20.

3.8. PSR J1010+15

A different type of pulse intensity modulation is observed in the case of RRATs, which emit pulses at a rate of a few per minute to a few per several hours. After enough pulses are observed, it is usually possible to determine a rotation period before (or without) any timing follow-up observations. Known RRATs appear to have longer periods on average than steadily emitting, non-recycled pulsars (McLaughlin et al. 2006; Deneva et al. 2009). However, when the detected number of pulses is small, the large intervals between successive pulses may cause the calculated period to be an integer multiple of the actual period.

As RRATs are a relatively new pulsar class, large surveys offer an opportunity to better understand their population statistics and the causes of their intermittency. In addition, since surveys covering high latitudes like AO327 look away from most of the ionized gas in our Galaxy, they are sensitive to extragalactic radio transients from potentially exotic sources such as merging and annihilating black holes. AO327 has discovered one RRAT to date, PSR J1010+15. This object has $\text{DM} = 42 \text{ pc cm}^{-3}$. Figure 10 shows the discovery observation, with two bright dispersed pulses detected ~ 18 s apart. There is no corresponding detection from the periodicity search branch of our pipeline.

4. FUTURE DIRECTIONS

We have described the current status and initial set of pulsar discoveries from the AO327 Arecibo drift survey. Phase 1 of the survey includes sky in the declination range -1° to 28° , excluding the region with $|b| < 5^\circ$, which is searched by the PALFA consortium at 1.4 GHz, a frequency better suited to the high-scattering and dispersion environment of the Galactic plane. To date we have completed observations of 52% of the Phase 1 region, processed 55% of the data, and discovered 24 new pulsars. We expect that most of the remaining Phase 1 region will be searched with the new PUPPI backend, which

Figure 10. Discovery observation of PSR J1010+15. The top panels show the number of pulses vs. S/N and DM, and S/N vs. DM for events with S/N > 5. At the bottom the same events are plotted vs. DM and time. Brighter pulses are represented by larger circles. The S/N peaks at DM $\sim$ 42.

provides finer time and frequency resolution, wider bandwidth, and therefore better sensitivity, especially to MSPs. Based on the current time AO327 time allocation of $\sim$ 400 hr per year, we expect to complete Phase 1 of the survey by 2017. Phase 2 of AO327 will target sky in the declination range 28° – 38° and we estimate that will take two additional years to complete.

Four of the 24 AO327 discoveries have $|b| < 10^\circ$, and one of these, PSR J1821+01, has $P = 0.034$ s. This pulsar is outside the PALFA search region within 5° of the Galactic plane. The fact that AO327 is finding fast pulsars at relatively low Galactic latitudes suggests that a non-drift search probing deeper parallel to the Galactic plane at $5^\circ \lesssim |b| \lesssim 10^\circ$ will likely be very fruitful and contribute to a more complete picture of the Galactic pulsar population.

When AO327 discovers a new pulsar, we routinely check for nearby sources from existing optical and high-energy catalogs, including the Second Fermi Source Catalog (2FGL; Nolan et al. 2012). Radio pulsar searches targeting unidentified 2FGL sources with pulsar-like gamma-ray spectra have shown that MSPs are common gamma-ray emitters (e.g., Ray et al. 2011). However, none of the pulsars in Table 2 have 2FGL counterparts. Conversely, once we obtain timing solutions from radio timing observations, we plan to use them to fold Fermi event data and check whether any of the AO327 pulsars exhibit gamma-ray pulsations, which may be too weak for the objects

to have been included in 2FGL. The ratio of radio-loud to radio-quiet gamma-ray pulsars is an especially good discriminator between gamma-ray pulsar population models, and radio surveys and multi-wavelength follow-up studies of new pulsars will make such distinction possible.

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